

Conservation of ζ from holography

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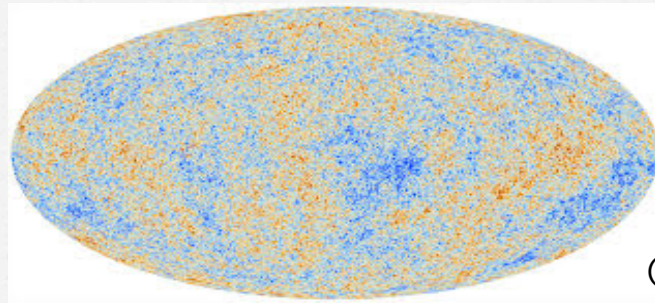
Y.U. & J.G. arXiv:1303.5997, JCAP 1307, 033

arXiv:1403.5497, JHEP 1406, 086

Y.U., J.G. & K.S. arXiv:1410.3290, JCAP 1501, ...

Inflation from holography

WMAP, PLANCK, ...



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Inflation is now quite compelling.

- Can we describe inflation holographically?

If YES, what's the prediction?

If NO, what's the obstacle?

Lyth bound

• Power spectrums $\Delta_{\zeta}^2 = \frac{1}{2M_{\text{pl}}^2} \frac{1}{\epsilon_*} \left(\frac{H_*}{2\pi} \right)^2$

$$\Delta_{\text{GW}}^2 = \frac{8}{M_{\text{pl}}^2} \left(\frac{H_*}{2\pi} \right)^2$$

$$r = \frac{\Delta_{\text{GW}}^2}{\Delta_{\zeta}^2} = 16\epsilon_*$$

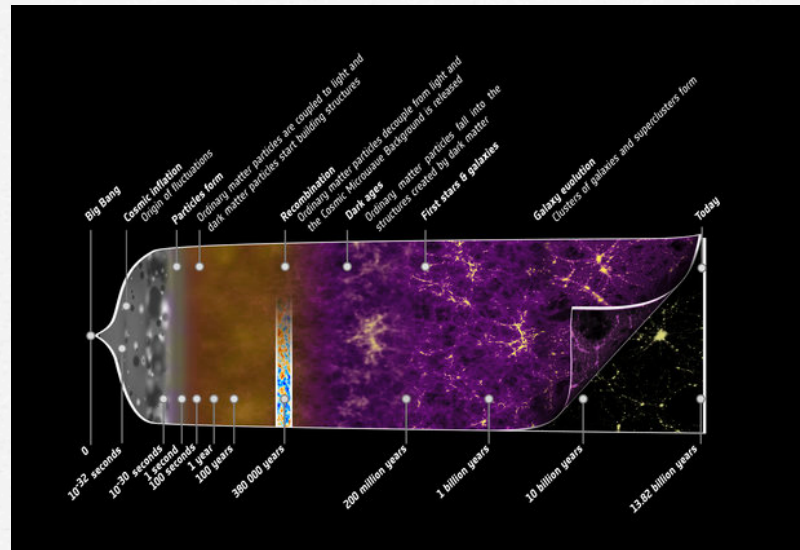
$$N_e = \int d \ln a = \int \frac{H}{\dot{\phi}} d\phi \sim \sqrt{8r}^{-1/2} \frac{\Delta\phi}{M_{\text{pl}}}$$

$$\longrightarrow \frac{\Delta\phi}{M_{\text{pl}}} \simeq \left(\frac{r}{0.1} \right)^{1/2}$$

$r \leq 0.12$ (Planck + BICEP 2015)

- Planck scale excursion is marginally allowed.

UV sensitivity of inflation



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energy

time

- Perturbation with controlled radiative corrections
 - High energy scale near Plank scale
 - Planck scale excursion is marginally allowed.

Outline of this talk

1. dS/CFT

2. Inflation/QFT

3. Boundary QFT

4. ζ correlators from boundary

skipped

Gauge/Gravity correspondence

- **Holographic principle**

't Hooft (92), Susskind (95)

Holographic principle suggests that a gravity theory should be related to a non-gravitational theory in one fewer dimension.

d-dim gauge theory \longleftrightarrow (d+1)-dim gravity theory
+ RG flow

- **Non-trivial duality**

Maldacena (97)

Boundary CFT

Bulk gravity

'tHooft coupling λ

$$\lambda = (r_0/l_s)^4$$

Curvature scale r_0

Strong coupling

$$\lambda \gg 1, r_0 \gg l_s$$

Weak coupling

Weak coupling

$$\lambda \ll 1, r_0 \ll l_s$$

Strong coupling

AdS/CFT as H-J formalism

d-dim gauge theory \longleftarrow (d+1)-dim gravity theory
+ RG flow

Recall Hamiltonian-Jacobi formalism....

using equation of motion for (d+1)-dim theory

$$\delta S \sim \mathcal{L} dz \Big|_{z=z_2}^{z=z_1}$$

(d+1)dim

$z=z_1$ holographic plane \rightarrow CFT

$z=z_2$ trivial B.C.

see also... holographic renormalization

S. Haro et al. (00), Skenderis (02),

Geometry of AdS and dS

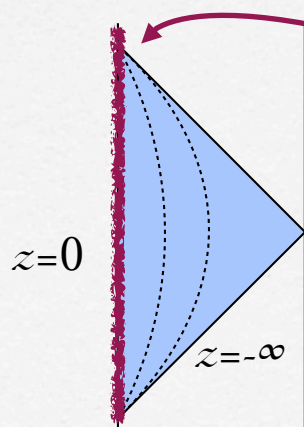
Anti de Sitter (AdS)

Vacuum with $\Lambda < 0$

in $\mathbb{R}^{2,3}$ $(-, -, +, +, +)$ $SO(2,3)$

$$-X_0^2 - X_1^2 + \sum_{a=2,3,4} X_a^2 = -A^2$$

$$ds^2 = l_{\text{AdS}}^2 \left(\frac{-dt^2 + dx^2 + dy^2 + dz^2}{z^2} \right)$$



Boundary

.....
z:const, \mathbb{R}^3



$$l_{\text{AdS}} \rightarrow i l_{\text{dS}}$$

$$z \rightarrow i\eta$$

$$t \rightarrow i\omega$$

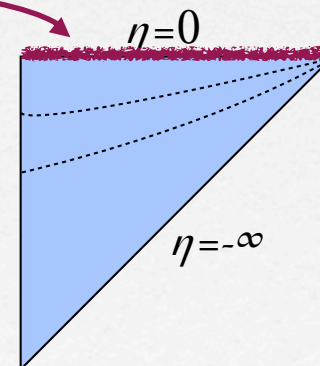
de Sitter (dS)

Vacuum with $\Lambda > 0$

in $\mathbb{R}^{1,4}$ $(-, +, +, +, +)$ $SO(1,4)$

$$-X_0^2 + X_1^2 + \sum_{a=2,3,4} X_a^2 = A^2$$

$$ds^2 = l_{\text{dS}}^2 \left(\frac{-d\eta^2 + dx^2 + dy^2 + d\omega^2}{\eta^2} \right)$$

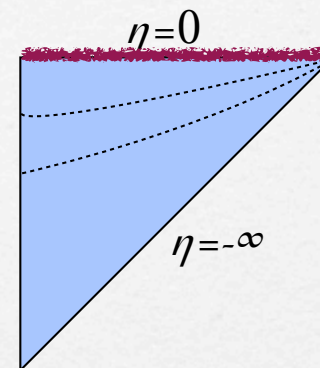


.....
 η :const, \mathbb{R}^3

dS/CFT

Strominger(01), Witten(01)

- CFT lives on the spacelike boundary at the future infinity of dS.



Maldacena(02)

- Wave function for bulk gravity
Probability distribution

$$\Psi_{\text{dS}}[g] = Z_{\text{CFT}}$$

$$P_{\text{dS}}[g] = |Z_{\text{CFT}}|^2$$

- Euclidean AdS $S_{\text{AdS ren}} = -\frac{1}{\kappa^2} \int \frac{d^3 k}{2\pi^3} \frac{1}{2} R_{\text{AdS}}^2 k^3 \phi_0(-k) \phi_0(k)$

$$\overrightarrow{Z \sim e^{-S}} \quad \langle O(k)O(-k) \rangle = -\frac{\delta^2 S_{\text{AdS ren}}}{\delta\phi_0(k)\delta\phi_0(-k)} = \frac{1}{2} \frac{R_{\text{AdS}}^2}{\kappa^2} k^3$$

- dS $\langle \phi(k)\phi(-k) \rangle = -\frac{1}{2\text{Re}\langle O(k)O(-k) \rangle} \Big|_{R_{\text{AdS}} = -iR_{\text{dS}}} = \frac{\kappa^2}{R_{\text{dS}}^2} \frac{1}{k^3}$

Analytic cont. connects dual boundaries of dS and AdS

Challenges of dS/CFT

- Holographic direction is time like.

Dual boundary theories to dS are non-unitary.

Good property?

- Poor understanding on analytic continuation.

Extendable to a non-perturbative example in $1/N$?

- Lack of a concrete example.

First concrete example of dS/CFT

Anninos, Hartman, & Strominger (11)

Vasiliev gravity in dS_4 \longleftrightarrow $Sp(N)$ CFT₃ living at \mathcal{I}^+

$$\Lambda \rightarrow -\Lambda$$

$$N \rightarrow -N$$

Outline of this talk

1. dS/CFT

2. Inflation/QFT

3. Holographic inflation (Simplest setup)

4. Conservation of ζ

Breaking symmetry

de Sitter space

4D hyperboloid:

$$ds_4^2 = \{\eta_{\mu\nu} X^\mu X^\nu = H^{-2}\}$$

in 5D flat spacetime

$SO(1,4)$
↔

CFT on R^3

- Poincare T.
- Dilatation
- Special C.T.

↓
Cosmological const. Λ
+ inflaton φ
Breaking dS sym.

Inflation

↓
CFT
+ φO (ex)mass
Breaking conf. sym.

Deformed CFT

Standard lore of inflation

4D bulk

Inflation

= dS + modulation

Given that....

- GR, $V(\varphi)$

- GR, $V(\varphi)$, $P(X=(\partial\varphi)^2)$

- $f(R)$, $V(\varphi)$ and so on

local QFT weakly coupled to gravity

→ $\varphi(t)$, $\langle \zeta \zeta \dots \zeta \rangle$,

Holographic inflation

4D bulk

Inflation
= dS + modulation

3D boundary

QFT
CFT+1 deformation



$$\Psi_{\text{bulk}}[\varphi] = Z_{\text{QFT}}[g]$$

$$Z_{\text{QFT}} = \int D\chi \exp \left[-S_{\text{CFT}} - \int \underline{g\mathcal{O}[\chi]} + \dots \right]$$

deformation

Necessary building blocks

- φ & g relation?
- t & μ relation?

Conservation of ζ

From cosmological perturbation

Single clock $\partial_t \zeta = O((k/aH)^2)$

wands et al. (00), Weinberg (03), Lyth et al (04),
Langlois & Vernizzi (05), ...

- Energy conservation $\nabla^\mu T^0_\mu = 0$
- Holds at full non-linear order

(ex) Single inflaton in Einstein gravity

$$\zeta'' + 2\frac{z'}{z}\zeta' - \cancel{\partial^2}\zeta = 0$$

Holographic inflation

4D bulk

Inflation
= dS + modulation

3D boundary

QFT
CFT+1 deformation



Conservation
of P_ζ

$$\Psi_{\text{bulk}}[\varphi] = Z_{\text{QFT}}[g]$$

$$Z_{\text{QFT}} = \int D\chi \exp \left[-S_{\text{CFT}} - \int g \mathcal{O}[\chi] \right]$$

{
- φ & g relation?
- t & μ relation?

Field redefinition

$$a(t) \propto \mu^C \quad C : \text{const} \quad \text{J.G.SY.U. (14)}$$

Outline of this talk

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$$Z_{\text{QFT}} = \int D\chi \exp \left[-S_{\text{CFT}} - \int g\mathcal{O}[\chi] \right]$$

4. Conservation of ζ

skipped

What's RG flow?

μ_p : Physical scale

μ : Cutoff scale

given that $\mu_p \ll \mu$



(ex) interaction

$$S = \prod_{i=1,2,3,4} \sum_{p_i < \mu} \varphi(p_i)$$

Physical quantity

$$F_{\text{phys}}(\mu_p; g(\mu), \mu) = F_{\text{phys}}(\mu_p; g(\mu'), \mu')$$

g : Physical constant



Rrenormalization group (RG) flow

Boundary QFT

Conformal perturbation theory (\sim Slow-roll expansion)

$$S_{\text{QFT}} = S_{\text{CFT}} + \delta S \quad \delta S = \int d^3x g \mathcal{O}[\chi] \quad (0 \leq g \ll 1)$$

\mathcal{O} : Boundary operator consists of χ

g : Dimensionless coupling

μ : Renormalization scale

- Correlators for CFT

$$\langle O(\mathbf{x})O(\mathbf{y}) \rangle_{\text{CFT}} = \frac{c}{|\mathbf{x} - \mathbf{y}|^{2\Delta}}$$

$$\langle O(\mathbf{x})O(\mathbf{y})O(\mathbf{z}) \rangle_{\text{CFT}} = \frac{C}{|\mathbf{x} - \mathbf{y}|^\Delta |\mathbf{y} - \mathbf{z}|^\Delta |\mathbf{z} - \mathbf{x}|^\Delta}$$

Beta function & FP

β function $\beta(\mu) \equiv \frac{dg(\mu)}{d \ln \mu}$

Klebanov et al. (11)

$$\beta(\mu) = \lambda g(\mu) + \frac{\tilde{C}}{2} g^2(\mu) + \mathcal{O}(g^3)$$

$$\tilde{C} \sim \frac{C}{c}$$

$$\lambda = \Delta - 3$$

Classical scaling

Quantum corrections

- Fixed point (FP) $\beta=0$

For $\tilde{C}/\lambda < 0$

Two FPs $g=0, -2\lambda/\tilde{C}$

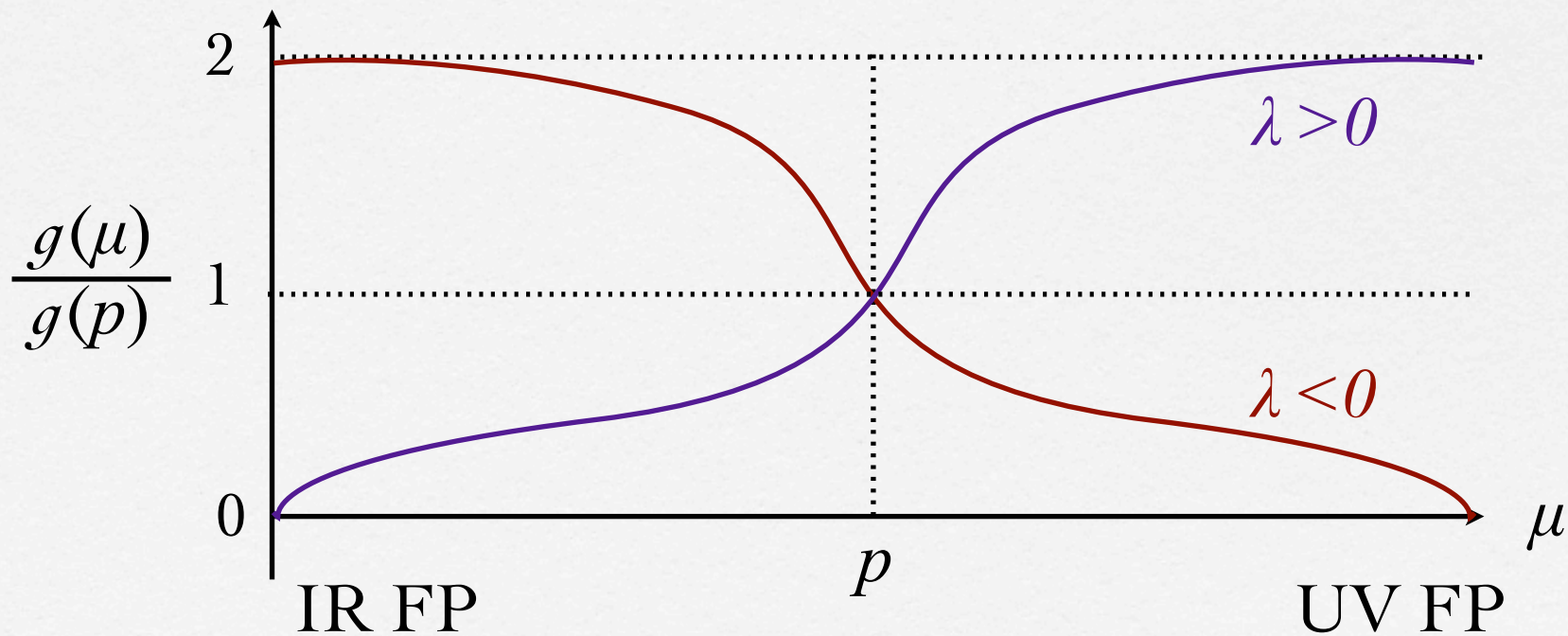
For $\tilde{C}/\lambda \geq 0$

One FP $g=0,$

Solving RG flow

$$\beta(\mu) = \lambda g(\mu) + \frac{\tilde{C}}{2} g^2(\mu) + \mathcal{O}(g^3) \quad \text{for } \tilde{C}/\lambda < 0$$

$$g(\mu) = \frac{2}{1 + \left(\frac{\mu}{p}\right)^\lambda} \left(\frac{\mu}{p}\right)^\lambda g(p) \quad g(p) \equiv -\frac{\lambda}{\tilde{C}}$$



Reconstruction of potential

KG equation

$$\ddot{\phi} + 3H\dot{\phi} + \frac{\partial V(\phi)}{\partial \phi} = 0$$

↓

$$\frac{d\phi}{d \ln a} = -\frac{2}{\kappa^2} \frac{1}{W(\phi)} \frac{\partial W(\phi)}{\partial \phi}$$

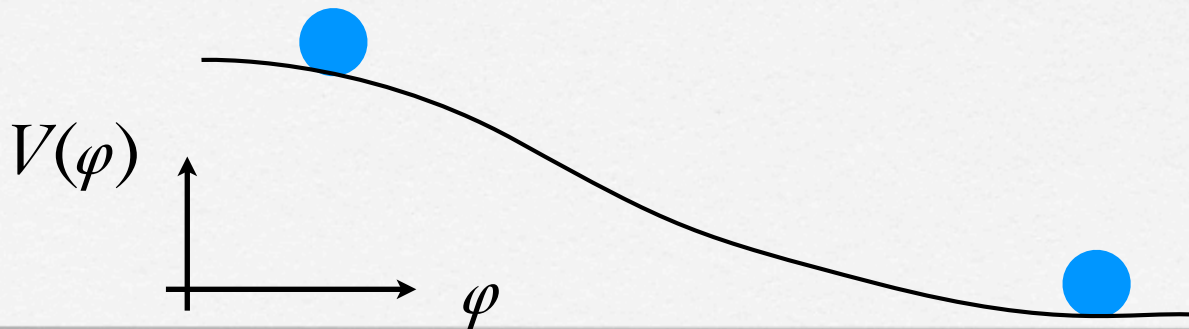
$$V(\phi) = \frac{8}{\kappa^2} \left[\frac{3}{2} W^2(\phi) - \frac{1}{\kappa^2} \left(\frac{\partial W(\phi)}{\partial \phi} \right)^2 \right]$$

RG equation

$$\frac{dg}{d \ln \mu} = \lambda g + \frac{\tilde{C}}{2} g^2 + O(g^3)$$

$\swarrow \searrow$

$$g(\mu, \mathbf{x}) = \kappa \phi(t(\mu), \mathbf{x})$$
$$d \ln a = p d \ln \mu$$



Correlators of \mathcal{O}

Expanding by correlators for CFT with cutoff

$$\begin{aligned} & \langle \mathcal{O}(x_1) \cdots \mathcal{O}(x_n) \rangle_\mu && \text{Bzowski et al. (12)} \\ & = \langle \mathcal{O}(x_1) \cdots \mathcal{O}(x_n) e^{-\int d^3x g \mathcal{O}} \rangle_{\mu, \text{CFT}} \end{aligned}$$

↓ integrating out $k > \mu$, changing μ , using OPE

$$Z^{-n/2}(\mu) \langle \mathcal{O}(x_1) \cdots \mathcal{O}(x_n) \rangle_{\mu, k < \mu} = Z^{-n/2}(\mu_0) \langle \mathcal{O}(x_1) \cdots \mathcal{O}(x_n) \rangle_{\mu_0, k < \mu_0}$$

Wave fn. renormalization

$$\sqrt{Z(\mu)} = \mu^{-\lambda} \left[1 + \left(\frac{\mu}{p} \right)^\lambda \right]^2 = 4p^{-\lambda} \frac{\beta(p)}{\beta(\mu)} \quad \text{J.G.Ψ.U. (14)}$$

Correlators

From boundary QFT to bulk gravity

$$\Psi_{\text{qdS}}[\delta\varphi] = Z_{\text{QFT}}[\delta\varphi] = e^{-W_{\text{QFT}}[\delta\varphi]} \quad P[\delta\varphi] = |\Psi_{\text{qdS}}[\delta\varphi]|^2$$

$$\langle \delta\phi(x_1) \cdots \delta\phi(x_n) \rangle = \int D\delta\phi P[\delta\phi] \delta\phi(x_1) \cdots \delta\phi(x_n)$$

* Distribution function $P[\delta\varphi] = e^{-\delta W[\delta\varphi]}$

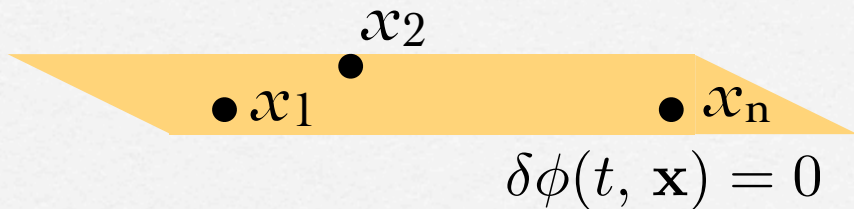
$$\delta W[\delta\phi] = \sum_{n=1}^{\infty} \int d^d \mathbf{x}_1 \cdots \int d^d \mathbf{x}_n W^{(n)}(\mathbf{x}_1, \cdots, \mathbf{x}_n) \delta\phi(\mathbf{x}_1) \cdots \delta\phi(\mathbf{x}_n)$$

$$W^{(n)}(\mathbf{x}_1, \cdots, \mathbf{x}_n) \equiv 2\text{Re} \left[\frac{\delta^n W_{\text{QFT}}[\delta\phi]}{\delta\phi(\mathbf{x}_1) \cdots \delta\phi(\mathbf{x}_n)} \Big|_{\delta\phi=0} \right]$$

ζ Correlators

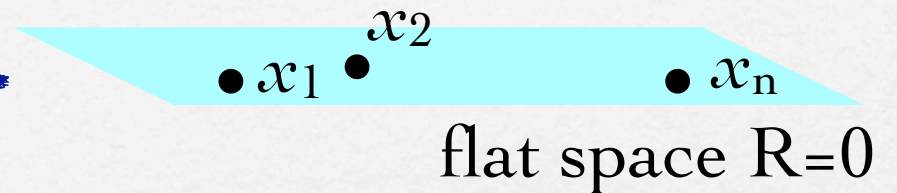
in cosmology

$$\langle \zeta(x_1) \zeta(x_2) \cdots \zeta(x_n) \rangle$$



in boundary QFT

$$\langle \mathcal{O}(x_1) \mathcal{O}(x_2) \cdots \mathcal{O}(x_n) \rangle$$



at large scales

$$\zeta = -\frac{H}{\dot{\phi}} \delta\phi + \frac{\varepsilon_2}{4} \left(\frac{H}{\dot{\phi}} \right)^2 \delta\phi^2 + \cdots$$

$$\varepsilon_1 \equiv \frac{1}{2} \frac{\dot{\phi}^2}{H^2}$$

$$\varepsilon_2 \equiv \frac{d \ln \varepsilon_1}{d \ln a}$$

Vertex function

$$W^{(n)}(x_1, \cdots, x_n) \equiv 2\text{Re} \left[\frac{\delta^n W_{\text{QFT}}[\zeta]}{\delta\zeta(x_1) \cdots \delta\zeta(x_n)} \Big|_{\zeta=0} \right] \longleftarrow \langle \mathcal{O}(x_1) \cdots \mathcal{O}(x_m) \rangle_\mu$$

$$\frac{\delta}{\delta\zeta} = \frac{\delta\phi}{\delta\zeta} \frac{\delta}{\delta\phi} \sim \frac{\delta\phi}{\delta\zeta} \mathcal{O}$$

Conservation of ζ

Power spectrum

J.G.εΥ.υ. (14)

$$\langle \zeta(\mathbf{x}_1) \zeta(\mathbf{x}_2) \rangle_{\text{conn}} = W^{(2)-1}(\mathbf{x}_1, \mathbf{x}_2)$$

$$B_1(\mu) = - \left. \frac{\partial \delta \phi}{\partial \zeta} \right|_{\zeta=0}$$

$$W^{(2)}(\mathbf{x}_1, \mathbf{x}_2) = -2 \text{Re} [B_1^2(\mu) \langle \mathcal{O}(\mathbf{x}_1) \mathcal{O}(\mathbf{x}_2) \rangle_\mu]$$

Conservation \longrightarrow $\frac{d}{d\mu} [B_1(\mu) \sqrt{Z(\mu)}] = 0$

Gauge transformation $B_1 = \frac{\dot{\phi}}{H} = \frac{d\phi}{d \ln a}$

RG flow $\sqrt{Z(\mu)} = 4p^{-\lambda} \frac{\beta(p)}{\beta(\mu)}$

Identification between t & μ

$$\ln(\mu/\mu_0) = C \ln(a/a_0)$$

C : const

$$C = 1 + O(\varepsilon)$$

Conserved Power spectrum

$$P(k) = -\frac{3}{8\pi} \frac{1}{c\beta^2(p)} \frac{1}{k^3} \left(\frac{k}{p}\right)^{-2\lambda} \left[1 + \left(\frac{k}{fp}\right)^\lambda\right]^4 \sim -\frac{1}{c\beta^2(k)}$$

cf Agrees with the result of Bzowski + (12) in $\mu \rightarrow \infty$

Remarks

1. Amplitude

$$\beta = \frac{dg}{d \ln \mu} \sim \frac{d(\phi/M_{pl})}{d \ln a} = \sqrt{2\varepsilon}$$

$$c \simeq (M_{pl}/H_{dS})^2 \quad \text{Strominger(01)}$$



$$\frac{1}{c\beta^2} \sim \frac{1}{\varepsilon} \left(\frac{H}{M_{pl}}\right)^2$$

Maldacena(02)

2. Spectral index

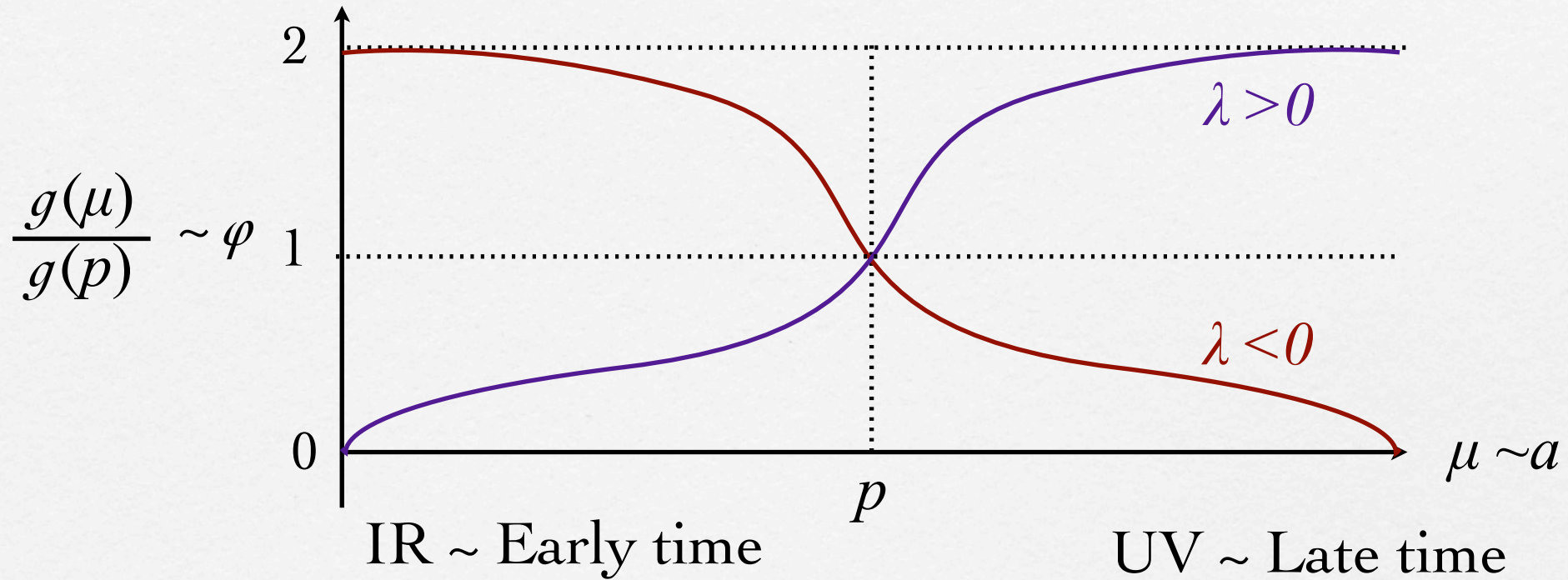
$$\text{For } k \gg fp \quad n_s - 1 = 2|\lambda|$$

Blue-tilted

$$\text{For } k \ll fp \quad n_s - 1 = -2|\lambda|$$

Red-tilted

Evolution of "inflaton"



N.B. $n_s - 1 = -6\epsilon + 2\eta$

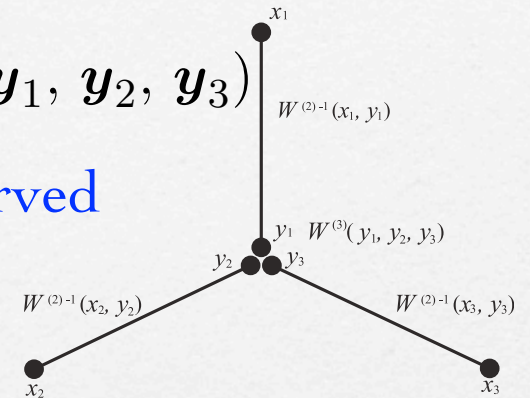
Conservation of bi-spectrum

$$\langle \zeta(\mathbf{x}_1) \zeta(\mathbf{x}_2) \zeta(\mathbf{x}_3) \rangle_{\text{conn}} = - \int \prod_{i=1}^3 d^d \mathbf{y}_i W^{(2)-1}(\mathbf{x}_i, \mathbf{y}_i) W^{(3)}(\mathbf{y}_1, \mathbf{y}_2, \mathbf{y}_3)$$

conserved if P_ζ conserved

$$W^{(3)}(\mathbf{y}_1, \mathbf{y}_2, \mathbf{y}_3) = -2\text{Re} \left[\beta^3(\mu) \langle \mathcal{O}(\mathbf{y}_1) \mathcal{O}(\mathbf{y}_2) \mathcal{O}(\mathbf{y}_3) \rangle \right.$$

$$\left. - \frac{d\beta(\mu)}{d \ln \mu} \beta(\mu) \delta(\mathbf{y}_1 - \mathbf{y}_2) \langle \mathcal{O}(\mathbf{y}_2) \mathcal{O}(\mathbf{y}_3) \rangle_\mu - (2 \text{ perms}) \right]$$



If the correlators of \mathcal{O} are given by

$$Z^{-n/2}(\mu) \langle \mathcal{O}(\mathbf{x}_1) \cdots \mathcal{O}(\mathbf{x}_n) \rangle_\mu = Z^{-n/2}(\mu_0) \langle \mathcal{O}(\mathbf{x}_1) \cdots \mathcal{O}(\mathbf{x}_n) \rangle_{\mu_0} \quad \sqrt{Z(\mu)} \propto 1/\beta(\mu)$$

Conservation requires

$$\frac{d \ln \beta}{d \ln \mu} = \text{const.} \quad \text{cosmologically} \quad \varepsilon_2 = \frac{d \ln \varepsilon_1}{d \ln a} = \text{const.}$$

Be more careful....

$$Z^{-n/2}(\mu) \langle \mathcal{O}(x_1) \cdots \mathcal{O}(x_n) \rangle_\mu = Z^{-n/2}(\mu_0) \langle \mathcal{O}(x_1) \cdots \mathcal{O}(x_n) \rangle_{\mu_0}$$

- RG solution does not apply to coincidence limit (CDL).
- We need regularization to compute CDL.

N.B. In CFT, symmetry argument does not specify the CDL.

$$S_{\text{QFT}} = S_{\text{CFT}} + \int d^d x g \mathcal{O}(x) + \int d^d x g_n \mathcal{O}^n(x)?$$

multi-trace operators

AdS/CFT

- Wilsonian RG, Bulk \rightarrow Bdry QFT with multi-trace op.
Heemskerk & Polchinski (10), Faulkner, Liu, & Rangamani (11)
- Bdry QFT with multi-trace op. \rightarrow GR (+ Λ)
S.S. Lee (13)

CONCLUSION

Holographic description of inflation scenario

- We computed the primordial spectrum holographically, and the result may apply to strong/weak gravity regimes (large N , arbitrary 'tHooft coupling).
- The conservation of ζ power spectrum determines t & μ relation as $a(t) \propto \mu^c$.
- A subtle issue on the conservation of bi-spectrum $\langle \zeta \zeta \zeta \rangle$
Yet, if we determine the CSL such that the consistency relation is fulfilled, the bispectrum is conserved.